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International Journal of HEAT and MASS TRANSFER

International Journal of Heat and Mass Transfer 47 (2004) 2487-2491

www.elsevier.com/locate/ijhmt

Technical Note

Full spectrum *k*-distribution correlations for CO₂ from the CDSD-1000 spectroscopic databank

Michael F. Modest *, Ranjan S. Mehta

Department of Mechanical Engineering, 137, Reber Building, Pennsylvania State University, University Park, PA 16802, USA Received 24 July 2003; received in revised form 22 November 2003

Keywords: Gas radiation; k-distributions; Engineering correlation

1. Introduction

Radiative heat transfer through molecular gases is a very important mode of energy transfer when dealing with combustion systems or modelling atmospheric processes. The most accurate radiative transfer results are obtained through line-by-line (LBL) calculations, which require, however, a prohibitively large computational effort. Consequently, a number of global and band models have been proposed. Presently, the most accurate global models are the full-spectrum k-distributions (FSK) by Modest and coworkers [1,2]. Earlier, somewhat less refined, related methods include the spectral-line-based weighted-sum-of-gray-gases [3-5] or SLW and absorption-distribution-function (ADF) [6,7] models. All of these methods are orders-of-magnitude more efficient than LBL calculations, and they all use FSK, which, in general, need to be calculated from highresolution databases, such as HITRAN [8], HITEMP [9] or-for CO₂-the new CDSD-1000 [10,11]. Both the SLW and ADF methods use simplified k-distributions reducing them to step-functions (gray gases), whereas the FSK method uses Gaussian quadrature to integrate over g-space resulting in better accuracy. Assembling such k-distributions is a rather tedious task and, therefore, to make simple engineering calculations feasible, Denison and Webb [12,13] have proposed several simple FSK correlations for CO2 and H2O, based on the outdated HITRAN92 [8] database, combined with some extrapolations of their own for high temperatures. Recently, Zhang and Modest [14] provided an updated

* Corresponding author. Fax: +1-814-863-8682.

E-mail address: mfm6@psu.edu (M.F. Modest).

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correlation for CO_2 based on the newer HITEMP database, which is purported to be accurate to at least 1000 K. Unfortunately, it appears that the HITEMP database shows some erroneous behavior for temperatures beyond 1200 K, as seen by comparison with experimental data [10,15]. The new CDSD-1000 databank [10,11], on the other hand, follows experimental data much more closely. Thus, it is the purpose of this note to provide a simple engineering correlation for FSK evaluated from the CDSD-1000 databank.

2. Mathematical formulation

The *k*-distributions are obtained by reordering the spectrally varying absorption coefficient into a monotonically increasing function with, in the case of FSK, the blackbody intensity (Planck function) as a weight factor. This has been described in detail in the original paper by Modest and Zhang [16].

The FSK is defined as

$$f(T_{\rm p}, T_{\rm g}, p, x; k) = \frac{1}{I_b} \int_0^\infty I_{b\eta}(T_{\rm p}) \delta(k - \kappa_\eta) \,\mathrm{d}\eta, \tag{1}$$

which is a function of temperature T_p through the black body intensity and of the state of the gas through the absorption coefficient $\kappa_{\eta}(T_g, p, x; \eta)$ where T_g is the gas temperature, p is the absolute pressure, x is the concentration and η is the wave number.

For RTE solution methods, the cumulative FSK is used, defined by

$$g(T_{\rm p}, T_{\rm g}, p, x; k) = \int_0^k f(T_{\rm p}, T_{\rm g}, p, x; k) \,\mathrm{d}k$$
$$= \frac{1}{I_b} \int_0^\infty I_{b\eta}(T_{\rm p}) H(k - \kappa_\eta) \,\mathrm{d}\eta, \tag{2}$$

where $H(k - \kappa_{\eta})$ is the Heaviside unit step-function. Physically, g is the Planck-function-weighted fraction of the spectrum over which the absorption coefficient $\kappa_{\eta} < k$.

The cumulative FSK $g(T_p, T_g, p, x; k)$, or rather its inverse $k(T_p, T_g, p, x; g)$ is used by today's sophisticated spectral solution methods, i.e., SLW [3,4,12,13,17], ADF [6,7,18], and FSK [1,2]. Determination of such cumulative k-distributions from high-resolution databases is a rather tedious and lengthy process and is, thus, not suitable for simple engineering calculations. With a small loss of accuracy, simple correlations may be employed instead, as first suggested by Denison and Webb using the outdated HITRAN92 [8] database. It is our aim here to present a similar but updated correlation for CO₂, employing the new CDSD-1000 databank [10,11]. Following Denison and Webb, we have

$$g(T_{\rm p}, T_{\rm g}, p, x; k) = \frac{1}{2} \tanh[P(T_{\rm p}, T_{\rm g}; k)] + \frac{1}{2}, \qquad (3)$$

where

$$P(T_{\rm g}, T_{\rm p}, k) = \sum_{l=0}^{3} \sum_{m=0}^{3} \sum_{n=0}^{3} d_{lmn} \left(\frac{T_{\rm g}}{1000K}\right)^{n} \\ \times \left(\frac{T_{p}}{1000K}\right)^{m} (\log_{10} k)^{l},$$
(4)

where $P(T_p, T_g; k) = P(T_p, T_g, p = 1 \text{ bar}, x = 0; k)$ with k in (cm bar)⁻¹, i.e., the correlation given here is valid only for the most common total pressure of $p \simeq 1$ bar, and, since CO₂ does not show strong self-broadening effects, the correlation is essentially independent of the CO₂ mole-fraction, which has thus been neglected. The errors that are introduced due to this approximation are small compared to errors of the correlated fit.

The correlation coefficients d_{lmn} listed in Table 1 were found by a least squares fit for FSK for CO₂ evaluated at 23 gas temperatures and 23 Planck function temperatures between 300 and 2500 K, with 10 orders of magnitude variation in k (cm bar)⁻¹ from $k = 10^{-8}$ (cm bar)⁻¹ to k = 200 (cm bar)⁻¹. Some representative k-distributions built from this correlation are plotted against the actual ones in Figs. 1 and 2. Unlike k-distributions calculated from HITRAN and HITEMP, those from CDSD-1000 show magnified undulations at



Fig. 1. Actual and approximate full-spectrum k-distribution functions for CO₂, calculated at a fixed Planck function temperature of 1000 K, with absorption coefficient taken at various temperatures.

Table 1 Coefficients for the full-spectrum k-distribution, d_{inm} , of carbon-dioxide for CDSD-1000

l	т	n				
		0	1	2	3	
0	0	1.85071	0.33373	0.62660	-0.12890	
	1	-0.20643	-2.57690	0.30090	-0.14090	
	2	0.27664	1.81420	-0.24728	0.10052	
	3	-0.37435E-01	-0.37762	0.53014E-01	-0.20836E-01	
1	0	0.67523	1.25760	0.67523E-01	-0.39669E-01	
	1	-0.70897	-3.07080	1.71150	-0.57694	
	2	0.48493	2.04603	-1.20220	0.40510	
	3	-0.98138E-01	-0.41928	0.24956	-0.84109E-01	
2	0	0.20690	0.28500	-0.48324E-01	0.15174E-01	
	1	-0.39473	-0.42333	0.54095	-0.23469	
	2	0.29020	0.21882	-0.34748	0.15857	
	3	-0.61998E-01	-0.38629E-01	0.68728E-01	-0.32314E-01	
3	0	0.38488E-01	0.18292E-01	-0.18958E-01	0.61307E-02	
	1	-0.41013E-01	-0.59115E-02	0.47118E-01	-0.23229E-01	
	2	0.37740E-01	-0.30114E-01	-0.11303E-01	0.12140E-01	
	3	-0.87906E-02	0.98357E-02	-0.53884E-04	-0.20604E-02	



Fig. 2. Actual and approximate full-spectrum k-distribution functions for CO₂, calculated at a fixed gas temperature of 2000 K, for different Planck function temperatures.

higher temperatures which are difficult to fit with a given number of coefficients. This is due to much more well defined rovibrational bands for CO_2 when modelled with CDSD-1000, in regions where HITEMP shows suspicious behavior from extrapolated lines in the band wings. However, as shown in sample heat transfer calculations below, the smoothening effect of the polynomial fit does not appear to seriously affect the accuracy of radiative calculations.

As a first check of the accuracy of the correlation, cumulative k-distributions are used to calculate the total emissivity of isothermal CO₂–N₂ mixtures. Fig. 3 shows total emissivities obtained directly from the CDSD and HITEMP databases (using LBL or, equivalently, FSK calculations), from the present correlation, the Denison and Webb correlation [12], and finally, from Leckner's correlation [19], which is based on fairly dated experimental data. As already noted by Zhang and Modest [14], HITEMP overpredicts emissivities at elevated temperatures, as compared with Leckner's model. Modest and Bharadwaj [15] have reported spurious absorption predicted by HITEMP at elevated temperatures in the band wings of the important 4.3 and 2.7 μ m bands of carbon dioxide. Emissivities from the CDSD-1000 databank and the present correlation are in better agreement with experimental data than HITEMP. Emissivities predicted by CDSD-1000 are generally lower than those measured by Leckner. This is an expected result since CDSD-1000 does not include all high-temperature lines. On the other hand, HITEMP emissivities become much too large at high temperatures due to the spurious lines, as discussed by Modest and Bharadwaj [15] and Tashkun et al. [10,11]. The old



Fig. 3. Total emissivities for different pressure path lengths from Leckner's correlation, direct calculations from HITEMP and CDSD-1000 databanks, and from the present correlation for CDSD.

correlation by Denison and Webb [3,4] does reasonably well, but underpredicts emissivities at high temperatures, due to missing hot lines. The disagreement at high pressure path lengths is assumed to be due to inaccuracy of the measured data. A similar trend, as shown in Fig. 4, is seen for the Planck-mean absorption coefficient obtained as outlined in [14], with HITEMP predicting much higher absorption at high temperatures than CDSD. The correlation does a reasonably good job approximating the Planck-mean absorption coefficient



Fig. 4. Pressure based Planck-mean absorption coefficient of CO_2 at different gas temperatures.

with a maximum error of about 10% at low temperatures.

Radiative heat loss from a simple one-dimensional layer between two parallel, cold, black walls is presented to further test the accuracy of the correlation, and is compared with "exact" (LBL or FSK) results obtained from the HITEMP and CDSD databanks. The medium is a N₂–CO₂ mixture at 1 bar total pressure and a CO₂ mole fraction of 10%, and the mixture temperature is varied from 300 to 2500 K. The flux exiting the walls is plotted in Fig. 5, showing a maximum error between the



Fig. 5. Radiative flux leaving a hot isothermal $90\%N_2-10\%CO_2$ mixture bounded by cold walls for varying gas temperatures.



Fig. 6. Radiative flux leaving from cold end of an isothermal $90\%N_2-10\%CO_2$ mixture at 1000 K between a hot and a cold, black wall for varying hot, black wall temperatures.

correlation and LBL calculations for CDSD of about 10%. As expected, at higher temperatures the heat-flux predicted by CDSD is less than that predicted by HITEMP, due to absence of the spurious lines that appear in the HITEMP spectrum at higher temperatures.

Finally, heat loss from a similar setup with one hot wall and one cold wall are predicted to test the accuracy of the correlation for transmission calculations. The mixture is kept at 1000 K with a 10% CO₂ mole fraction and hot wall temperature varying from 300 to 2500 K. Shown in Fig. 6 are the fluxes leaving the cold end. The error between LBL results obtained from the CDSD databank and from the present correlation are always less than 10%.

3. Conclusion

A simple correlation for full-spectrum k-distributions of N_2 -CO₂ mixtures was developed, based on the new spectroscopic databank CDSD-1000. The correlation was tested through typical one-dimensional problems and found to achieve reasonable accuracy as compared to FSK obtained directly from the database. Thus the correlation allows simple, yet accurate determination of radiative fluxes in CO₂ mixtures, without need of tedious calculations employing a high-resolution databank.

Acknowledgements

The authors gratefully acknowledge the financial support of the National Science Foundation under contract CTS-0112423.

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